Weak evolution of the mass–metallicity relation at cosmic dawn in the FirstLight simulations

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ABSTRACT

Little is known about the mass-metallicity relation (MZR) in galaxies at cosmic dawn. Studying the first appearance of the MZR is one of the keys to understand the formation and evolution of the first galaxies. In order to lay the groundwork for upcoming observational campaigns, we analyse 290 galaxies in haloes spanning $M_h = 10^9 - 10^{11} M_{\odot}$ selected from the FirstLight (FL) cosmological zoom simulations to predict the MZR at z = 5-8. Over this interval, the metallicity of FL galaxies with stellar mass $M_* = 10^8 M_{\odot}$ declines by $\leq 0.2 \text{ dex}$. This contrasts with the observed tendency for metallicities to increase at lower redshifts, and reflects weakly evolving or even increasing gas fractions. We assess the use of the R3 strong-line diagnostic as a metallicity indicator, finding that it is informative for $12 + \log (O/H) < 8$ but saturates to $R3 \approx 3$ at higher metallicities owing to a cancellation between enrichment and spectral softening. None the less, campaigns with *JWST* should be able to detect a clear trend between R3 and stellar mass for $M_* > 10^{7.5} M_{\odot}$. We caution that, at fixed metallicity, galaxies with higher specific star formation show higher R3 owing to their more intense radiation fields, indicating a potential for selection biases.

Key words: galaxies: evolution - galaxies: formation - galaxies: high-redshift.

1 INTRODUCTION

Little is known about the abundances of elements heavier than helium in the primeval galaxies of the early Universe. These elements are thought to be produced in the first supernovae explosions (Bromm & Yoshida 2011). Therefore, the generation of heavy elements was quickly linked to star formation processes and galaxy growth, driven by gas inflows, outflows, and merging. In fact, the evolution of the metallicity, the content of heavy elements relative to hydrogen and helium, gives strong insights about these processes (Maiolino & Mannucci 2019).

The galaxy average gas-phase metallicity shows a strong scaling relation with the galaxy stellar mass. The mass-metallicity relation (MZR) has been observed from the local Universe (z = 0) to cosmic noon ($z \simeq 3.5$; see Maiolino & Mannucci 2019 for a review of observational efforts). There is a consensus that the MZR evolves at z < 3.5 in the sense that, at fixed stellar mass, the metallicity declines with redshift. Does this trend continue to even higher

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redshifts? There are no observational estimates of metallicities at cosmic dawn ($z \ge 5$) with the exception of GRB afterglows (Berger et al. 2007). Measurements at lower redshifts ($z \le 3$) are mainly based on strong emission lines, such as [OIII] λ 5007. At cosmic dawn, these optical lines are redshifted out of the spectral range of current spectrographs. The James Webb Space Telescope (*JWST*) will soon open a new window into this range, allowing the MZR to traced to much earlier times. Therefore, it is now the best time to make theoretical predictions about the evolution of the MZR at cosmic dawn.

Cosmological simulations have become very advanced tools to study the evolution of the MZR (Finlator & Davé 2008; De Rossi et al. 2017; Torrey et al. 2018; Davé et al. 2019). However, there are only a few cosmological simulations that provide results at cosmic dawn (Barrow et al. 2017; Ceverino, Glover & Klessen 2017; Finlator et al. 2018; Rosdahl et al. 2018; Katz et al. 2019; Ma et al. 2019; Pallottini et al. 2019; Torrey et al. 2019). Most of these works achieve a high resolution using the 'zoom-in' technique, which concentrates all computational power on a few select haloes.

The main advantage of zoom simulations over full-box simulations is that they treat the baryon cycle with much more realism. For example, they model both the emergence of star formation-driven outflows and the small-scale interactions between outflows, the host galaxy's interstellar medium, and the circumgalactic medium more accurately due to their higher resolution and more detailed models of star formation and feedback. As these processes of cooling, star formation, ejection, mixing, and re-accretion are the processes that govern the baryon cycle and the metal enrichment within galaxies, it is important to compare predictions emerging from both simulation techniques.

The traditional drawback to zoom-in simulations is that they do not reveal the ensemble statistical properties of galaxy populations such as the MZR and its scatter. Consequently, most of works focus on the formation of few galaxies. The FirstLight (FL, Paper I Ceverino et al. 2017) data base of ~300 zoom-in simulations overcomes this limitation. Moreover, FL galaxies agree well with a variety of observations at z = 5-8, such as the UV luminosity function or the stellar mass function (Paper I), the SFR– M_* relation (Paper II, Ceverino, Klessen & Glover 2018) and M_* –Luminosity relations (Paper III, Ceverino, Klessen & Glover 2019). Thus, the FL data base is the ideal laboratory to study the predicted MZR at $z \ge 5$.

This paper has two well-defined goals. After the description of simulations (Section 2), we characterize the MZR at different redshifts at cosmic dawn (Section 3.1) and study its evolution (Section 3.2). Then, we provide metal-sensitive observables (Section 3.3) that can be used in the calibration of future measurements of metallicity at cosmic dawn. Section 4 ends with the summary and discussion.

2 SIMULATIONS

This paper uses a complete mass-selected subsample from the FL data base of simulated galaxies, described fully in Paper I. The subsample consists of 290 haloes with a maximum circular velocity (V_{max}) between 50 and 250 km s⁻¹, selected at z = 5. The haloes cover a mass range between a few times 10⁹ and 10¹¹ M_☉. This range excludes more massive and rare haloes with number densities lower than $\sim 3 \times 10^{-4} (h^{-1} \text{ Mpc})^{-3}$, as well as small haloes in which galaxy formation is extremely inefficient.

The target haloes are initially selected using low-resolution *N*body only simulations of two cosmological boxes with sizes 10 and 20 h^{-1} Mpc, assuming WMAP5 cosmology with $\Omega_m = 0.27$, $\Omega_b = 0.045$, h = 0.7, and $\sigma_8 = 0.82$ (Komatsu et al. 2009). We select all distinct haloes with V_{max} at z = 5 greater than a specified threshold, log $V_{\text{cut}} = 1.7$ in the 10 h^{-1} Mpc box and log $V_{\text{cut}} = 2.0$ in the 20 h^{-1} Mpc box. Initial conditions for the selected haloes with much higher resolution are then generated using a standard zoomin technique (Klypin, Trujillo-Gomez & Primack 2011). The DM particle mass resolution is $m_{\text{DM}} = 10^4 \text{ M}_{\odot}$. The minimum mass of star particles is 100 M_{\odot} . The maximum spatial resolution is always between 8.7 and 17 proper pc (a comoving resolution of 109 pc after z = 11).

The simulations are performed with the ART code (Kravtsov, Klypin & Khokhlov 1997; Kravtsov 2003; Ceverino & Klypin 2009; Ceverino et al. 2014, Paper I) that accurately follows the evolution of a gravitating *N*-body system and Eulerian gas dynamics using an adaptive mesh refinement (AMR) approach. Besides gravity and hydrodynamics, the code incorporates many of the astrophysical processes relevant for galaxy formation. These processes, represented via subgrid physical prescriptions, include gas cooling due to atomic hydrogen and helium, metal and molecular hydrogen cooling, photoionization heating by a constant cosmological UV

background with partial self-shielding, star formation, and feedback (thermal + kinetic + radiative), as described in Paper I. The simulations track metals released from SNe-Ia and from SNe-II, using supernovae yields that approximate the results from Woosley & Weaver (1995). These values are given for gas cells and star particles as described in Kravtsov (2003).

We assume that the unresolved nebular region around each star particle shares the same mass ratio of metals produced in SN II explosions as in the star particle (Paper III). The galaxy metallicity is defined as the mass-weighted average nebular metallicity including all star particles younger than 100 Myr. Using the supernovae yields included in the simulation (Woosley & Weaver 1995), our definition of solar metallicity corresponds to $Z_{\odot} = 0.02$ and $\log(O/H) + 12 = 8.9$. This normalization differs from other simulation works, such as Torrey et al. (2019), which assume a slightly lower value of 8.6. This systematic difference in normalization is comparable to the ~factor-of-two uncertainties in observational metallicity calibrations (Kewley & Ellison 2008). It does not impact our current study, which focuses on the relative evolution of the MZR at cosmic dawn.

The luminosities of metal-sensitive emission lines are extracted from the publicly available SEDs described in Paper III. In summary, the SEDs of the simulated galaxies are computed using publicly available tables from the Binary Population and Spectral Synthesis (BPASS) model (Eldridge et al. 2017) including nebular emission (Xiao, Stanway & Eldridge 2018). We combine all individual SEDs coming from all star particles within each galaxy. The nebular emission originates in regions around star particles younger than 100 Myr.

As described in Xiao et al. (2018), we compute the line luminosities of a single nebular region around a star particle by assuming a constant nebular density of $n_{\rm H} = 300 \,{\rm cm}^{-3}$, because the simulations do not resolve the nebular regions around young stars where most of the nebular light is emitted. This is the value normally used in the literature (Steidel et al. 2016). Denser H II regions, $n_{\rm H} = 1000 \,{\rm cm}^{-3}$, give similar results for the SSP metallicities considered in this paper (Xiao et al. 2018). Next, we measure the ionization parameter at the Strömgren radius. The Strömgren radius for each star particle is calculated using the assumed nebular density and the properties of the stellar population. Finally, we use a publicly available grid of Cloudy models to derive the luminosities of the most prominent optical and UV emission lines. More details can be found in Xiao et al. (2018).

Radiative transfer effects and dust attenuation are not included in the analysis. These could affect the rest-frame optical line luminosities of the most massive galaxies, $M_* > 10^8 \text{ M}_{\odot}$. Paper III has shown that the computed dust-free UV slope (β) is consistent with observations (Bouwens et al. 2014) for galaxies fainter than $M_{\text{UV}} \simeq -19$. In other words, available observations are consistent with the assumption that dust obscuration has a minor impact on the luminosity ratios of rest-frame optical lines for most of the FL sample.

3 RESULTS

3.1 The mass-metallicity relation

In Fig. 1, we show that a tight MZR is predicted to exist throughout z = 5-8. Moreover, it evolves very little with redshift. We do not see a change in the slope, as reported by Torrey et al. (2019) at $M_* \simeq 10^9 \,\mathrm{M_{\odot}}$. The shallower slope in FL results from the minimum wind velocity imposed in the recent Illustris model (Pillepich et al. 2018).



Figure 1. Metallicity as a function of the stellar mass from redshift z = 5-8. The colour bar represents the gas fraction. The red lines mark the linear fit. There is weak evolution of this scaling relation at cosmic dawn.

This assumption reduces the mass loading factor (that is, the gas outflow rate in units of SFR) in comparison to what occurs without such a velocity floor (Torrey et al. 2019). A reduced loading factor naturally increases the galaxy metallicity at a fixed stellar mass because more metals are retained in the ISM rather than expelled (Lilly et al. 2013; Finlator 2017).

Observations indicate that the scatter in the MZR correlates with other galaxy properties including gas fraction and SFR (Maiolino & Mannucci 2019). Previous cosmological full-box simulations have reported this correlation at lower redshifts, $z \le 4$ (De Rossi et al. 2017; Davé et al. 2019; Torrey et al. 2019). In order to explore whether this is expected even at $z \ge 6$, we use colours to indicate the gas fraction F_{gas} , defined as $F_{\text{gas}} \equiv M_{\text{gas}}/(M_{\text{gas}} + M_*)$ in Fig. 1. Indeed, it does seem that a residual correlation between Z and F_{gas} contributes significantly to the MZR scatter. For example, galaxies with higher-than-average gas fractions, $F_{\rm gas}\simeq 0.7$, have preferentially lower metallicities, at a fixed stellar mass of $10^8\,M_\odot.$ In the same mass bin, galaxies with lower-than-average gas fractions, $F_{\rm gas}$ \simeq 0.3, have metallicities higher by 0.2 dex than average. As shown in Paper II, a higher gas fraction correlates with a higher star formation rate for a given mass and redshift. Therefore, the FL simulations predict that the fundamental mass-metallicity relation (Ellison et al. 2008; Maiolino & Mannucci 2019) has already emerged by cosmic dawn.

3.2 Weak evolution of the MZR

Fig. 2 (blue circles) shows the average metallicity and its scatter (vertical error bars) within a narrow mass bin centred at $M_* = 10^8 \,\mathrm{M}_{\odot}$ at different redshifts across cosmic dawn. These average values agree well with the linear fits used in Fig. 1 (yellow diamonds). They show weak evolution of the MZR at these redshifts. In fact, there is even a hint that metallicity declines by ~0.15 dex from z = 8-5, although the decline is similar to the intrinsic scatter of the relation.

We can show that this weak evolution is driven by weak evolution in F_{gas} by invoking the *effective yield* (Garnett 2002; Dalcanton 2007). If a galaxy is a closed-box system, then the gas metallicity

$$Z_{\rm g} \equiv M_{\rm metals}/M_{\rm gas} \tag{1}$$

is a function of F_{gas} and the intrinsic stellar yield, y = 0.02, which expresses the ratio of the mass of new metals released into the ISM to the mass in long-lived stars. In the more general case where inflows and outflows are permitted, y_{eff} quantifies their net impact on a galaxy's chemical content

$$Z_{\rm g} = -y_{\rm eff} \ln(F_{\rm gas}). \tag{2}$$

Under the assumption that the net impact of inflows and outflows is constant throughout our redshift range and stellar mass bin, y_{eff} is



Figure 2. The evolution of metallicity at fixed stellar mass, $M_* = 10^8 \text{ M}_{\odot}$. Blue circles and yellow squares indicate FL predictions while coloured curves indicate predictions from other simulations (dashed lines indicate an extrapolation in redshift). Red triangles show the predicted evolution if the effective yield is fixed to $y_{\text{eff}} = 0.002$ and $F_{\text{gas}}(z, M_* = 10^8 \text{ M}_{\odot})$ is adopted directly from our simulations (Section 3.2). Note that yellow squares and red triangles are offset slightly in redshift for clarity.

constant and the MZR's evolution in time at a fixed stellar mass is driven by evolution in F_{gas} . The average gas fractions increase from $F_{\rm gas} \simeq 0.45$ at z = 8 to $F_{\rm gas} \simeq 0.55$ at z = 6. In order to verify that this suffices to drive the predicted metallicity evolution, we use red triangles in Fig. 1 to plot the quantity $-0.002\ln(F_{gas})$. This model closely tracks the FL simulations throughout $z = 8 \rightarrow 6$, supporting the idea that slowly increasing gas fractions indeed dominate early metallicity evolution within the full simulation. From $z = 6 \rightarrow 5$, the agreement is weaker, although the constant effective yield model still agrees with the simulation to within the simulated scatter. For context, we note that the assumed value $y_{eff} = 0.002$ is similar to the values found in local galaxies of similar mass (Tremonti et al. 2004). The effective yield is 10 times smaller than the intrinsic yield, indicating that the interstellar media of reionizationepoch galaxies were highly dynamical environments characterized by a vigorous interplay between metal-rich outflows and pristine inflows.

The slowly increasing gas fractions that arise in our simulations predict that, prior to $z \simeq 5$, galaxies build up their gas reservoirs because the rate at which star formation and outflows process their ISM lags the gas accretion rate. In this regime, dilution from pristine inflows dominates over enrichment, leading to constant or even slowly decreasing gas-phase metallicities (see also Wu et al. 2019).

As the MZR at z > 4 remains unconstrained observationally, we focus on a comparison with predictions from other theoretical models. Ma et al. (2016b) show results up to z = 6 and Torrey et al. (2019) extend it to z = 7. In order to facilitate comparison with FL predictions, we extrapolate their fitting relations to our range of redshifts. In doing so, we normalize their predictions to the same solar abundance that we assume (see Section 2). Our metallicity values lie in between these predictions. The values are close, within ~0.5 dex. This is encouraging because these simulations use slightly different definitions of metallicity but they claim to reproduce the same observed values at lower redshifts, $z \leq 2$.

However, there are some important differences in the evolution of the MZR with respect to these previous works. In Torrey et al. (2019), we can see that the metallicity strongly decreases with redshift at all times. This evolution is driven by the artificial velocity floor imposed on the outflows, as described above. Indeed, their massive galaxies, $M_* \ge 10^{9.5} \,\mathrm{M_{\odot}}$, start to show a flattening in their metallicity evolution with redshift at $z \simeq 5$. This is because more massive galaxies have outflows with higher velocities and they are not affected by this velocity floor.

The simulations by Finlator et al. (2018) show no metallicity evolution from z = 3-6, roughly consistent with our findings at higher redshifts. Although their wind model is similar to Torrey et al. (2019), it does not include the velocity floor. Most probably, this is why their metallicity does not decrease with redshift as in previous works. This highlights the importance of galactic outflows in the evolution of the metal content in galaxies.

The simulations by Ma et al. (2016a) predict a flattening of the metallicity evolution with increasing redshift, roughly consistent with the weak evolution in FL. However, the predicted metallicities are systematically lower by 0.3 dex. It is likely that the outflows in these simulations are far more efficient in removing metals from small galaxies. Indeed, Agertz et al. (2020) have shown that this model struggles to reproduce the observed plateau of stellar abundances in the faintest dwarfs in the local volume.

These comparisons highlight the importance of metallicity as a sensitive test of feedback models. Two questions motivated by this discussion include. First, when are galaxies predicted to graduate from building up their gas reservoirs to exhausting them? The epoch at which F_{gas} peaks may manifest observationally via a minimum in the MZR normalization. Secondly, how can upcoming measurements of F_{gas} and Z_{gas} be used to identify this transition? Over the next decade, observations of galaxies and their circumgalactic media at cosmic dawn will distinguish between these models, illuminating how the MZR first emerged.

3.3 Metallicity indicators

We now assess the strong rest-frame optical emission line diagnostics that will soon be used to test the predictions in the previous sections. These lines are widely used to infer the metallicities of low-redshift galaxies because they are bright and easily accessible using ground-based measurements. At reionization-epoch redshifts, however, only *JWST* can open up rest-frame optical diagnostics of galaxy evolution.

From the published mock optical spectra of the galaxies in the FL data base (Paper III), we measure the luminosities of [OIII] λ 5007, OII (\equiv [OII] λ 3727 + [OII] λ 3729), and H β . From these fluxes, we compute the metallicity-sensitive *R2* and *R3* indices following Maiolino & Mannucci (2019). As oxygen is the most abundant heavy element, we can compare them with the intrinsic gas-phase metallicities discussed in the previous section. In practice, we have found that the contribution of *R2* to *R23* (which is the sum or *R2* and *R3*) is very small. We therefore focus on *R3*, with the understanding that qualitative results apply also to *R23*.

In Fig. 3, we plot R3 versus the galaxy metallicity. For 12 + log (O/H) < 8, R3 increases strongly with gas-phase metallicity. At higher metallicities, it saturates. This behaviour reflects a competition between enrichment and ionization. For 12 + log (O/H) < 8, a slight increase in metallicity boosts the oxygen abundance more significantly than it softens the radiation field, yielding an overall increase to R3. At higher metallicities, the two effects largely cancel. As a result, R3 is a useful probe of metallicity only at low metallicities.

The scatter of this relation is mostly driven by the sSFR (sSFR \equiv SFR/ M_*), which correlates with the ionization parameter (Paper III). At a given metallicity, galaxies with higher sSFR have



Figure 3. *R3* as a function of the metallicity at redshift z = 6. The colour bar represents sSFR. After a maximum value at around $Z = 0.1 Z_{\odot}$, *R3* strongly decreases at lower metallicities. For a given metallicity bin, *R3* increases with sSFR.



Figure 4. *R3* as a function of the stellar mass at redshift z = 6. The colour bar represents the metallicity in oxygen units. There is drop in *R3* which corresponds to a metallicity around 7.5 and a stellar mass around $10^{7.5} M_{\odot}$.

systematically higher *R3* values. These galaxies are usually brighter and easier to observe. This could introduce a bias towards high *R3* values in the observations of the faintest galaxies.

In Fig. 4, we plot *R3* versus the galaxy stellar mass. Here, we see a clear trend of increasing *R3* with mass. The trend is mostly driven by the galaxy metallicity. There is a drop in *R3* at $M_* \simeq 10^7 \,\mathrm{M_{\odot}}$, driven by the low oxygen abundance, as in Fig. 3. JWST will only just be capable of observing such low-mass galaxies in blank fields. Ultradeep fields with a limiting magnitude of m = 31 in the rest-frame UV would be mass completed for $M_* \ge 10^{7.5} \,\mathrm{M_{\odot}}$ at z = 6 (Paper III). JWST samples of smaller galaxies will be biased towards higher sSFRs, which could prevent detection of the predicted drop in *R3* towards lower masses.

4 DISCUSSION AND SUMMARY

We have used the FL data base of cosmological zoom-in simulations to study the mass-metallicity relation at cosmic dawn. The main results can be summarized as follows: (i) The MZR, well-known at low redshifts, is predicted to exist at high redshifts (z = 5-8).

(ii) The scatter of the MZR is driven by the gas fraction.

(iii) The average metallicity at a fixed stellar mass evolves weakly – and may even decrease – during the interval $z = 8 \rightarrow 5$. This evolution is driven by slowly increasing gas fractions during the epoch of gas reservoir buildup.

(iv) The emission lines ratios *R3* and *R23* correlate relatively well with both metallicity and stellar mass. However, a secondary dependence on sSFR could introduce observational biases, particularly in samples that are not mass selected.

(v) There is an abrupt decrease of R3 at low metallicities and low masses driven by low metal abundances.

Galaxy spectra from JWST will open a new window to the restframe optical at cosmic dawn (Álvarez-Márquez et al. 2019). New measurements of emission lines ratios at these high redshifts will give us the first direct determination of galaxy metallicities in the early Universe. However, a lot of work on calibration of these metal tracers will be needed. This work provides the first attempt to look at these metallicity tracers from cosmological simulations. Still, systematic uncertainties remain. The choice of different line-emission models could affect the calibration. The choice of stellar evolution model (in our case, BPASS) is another source of uncertainty. Future works will address the systematics associated with these decisions and the effect in the overall normalization. However, the reported evolutionary trends are robust to choices that only affect the MZR's normalization.

Other caveats are intrinsically related to the FL simulations. First, they do not follow individual elements, like oxygen, carbon, or nitrogen. Therefore, we rely on published supernovae yields to estimate the oxygen abundances from the total amount of metals produced in core-collapsed supernovae. This is a good assumption for oxygen because it is mostly generated in these supernovae. However, other important elements, like carbon or nitrogen, have also secondary production channels, like AGB winds, which are not included. In addition, radiative transfer effects are only considered at post-processing on unresolved scales. The structure and dynamics of nebular regions on very small scales may also affect O III luminosities (Pellegrini et al. 2019). Dust is also not included, although its effect on R3 is small. Finally, bigger cosmological volumes will be needed to address the metal content in more massive galaxies and/or higher redshifts.

There is still plenty of exciting work to do, things to try and to understand, with respect to the mechanisms of metal production in the early Universe, or more precisely in primeval galaxies. Future works include following the metallicity content of galaxies in the same mass range at z = 5-8 by computing the 3D metal distributions in time intervals of 10 Myr. As a second step, we will follow the evolution of M_{metals} and M_{gas} for different galaxies at similar masses but different redshifts. We will check if their evolution could explain the mild increase of metallicity we get with redshift. These followup works will get new insights about the origin of the elements in the early Universe.

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REFERENCES

- Agertz O. et al., 2020, MNRAS, 491, 1656
- Álvarez-Márquez J. et al., 2019, A&A, 629, A9
- Barrow K. S. S., Wise J. H., Norman M. L., O'Shea B. W., Xu H., 2017, MNRAS, 469, 4863
- Berger E. et al., 2007, ApJ, 665, 102
- Bouwens R. J. et al., 2014, ApJ, 793, 115
- Bromm V., Yoshida N., 2011, ARA&A, 49, 373
- Ceverino D., Klypin A., 2009, ApJ, 695, 292
- Ceverino D., Klypin A., Klimek E. S., Trujillo-Gomez S., Churchill C. W., Primack J., Dekel A., 2014, MNRAS, 442, 1545
- Ceverino D., Glover S. C. O., Klessen R. S., 2017, MNRAS, 470, 2791
- Ceverino D., Klessen R. S., Glover S. C. O., 2018, MNRAS, 480, 4842
- Ceverino D., Klessen R. S., Glover S. C. O., 2019, MNRAS, 484, 1366
- Dalcanton J. J., 2007, ApJ, 658, 941
- Davé R., Anglés-Alcázar D., Narayanan D., Li Q., Rafieferantsoa M. H., Appleby S., 2019, MNRAS, 486, 2827
- De Rossi M. E., Bower R. G., Font A. S., Schaye J., Theuns T., 2017, MNRAS, 472, 3354
- Eldridge J. J., Stanway E. R., Xiao L., McClelland L. A. S., Taylor G., Ng M., Greis S. M. L., Bray J. C., 2017, Publ. Astron. Soc. Aust., 34, e058
- Ellison S. L., Patton D. R., Simard L., McConnachie A. W., 2008, ApJ, 672, L107
- Finlator K., 2017, in Fox A., Davé R., eds, Astrophysics and Space Science Library, Vol. 430, Gas Accretion on to Galaxies. Springer-Verlag, Berlin, p. 221
- Finlator K., Davé R., 2008, MNRAS, 385, 2181
- Finlator K., Keating L., Oppenheimer B. D., Davé R., Zackrisson E., 2018, MNRAS, 480, 2628
- Garnett D. R., 2002, ApJ, 581, 1019

- Katz H. et al., 2019, MNRAS, 487, 5902
- Kewley L. J., Ellison S. L., 2008, ApJ, 681, 1183
- Klypin A. A., Trujillo-Gomez S., Primack J., 2011, ApJ, 740, 102
- Komatsu E. et al., 2009, ApJS, 180, 330
- Kravtsov A. V., 2003, ApJ, 590, L1
- Kravtsov A. V., Klypin A. A., Khokhlov A. M., 1997, ApJS, 111, 73
- Lilly S. J., Carollo C. M., Pipino A., Renzini A., Peng Y., 2013, ApJ, 772, 119
- Maiolino R., Mannucci F., 2019, A&AR, 27, 3
- Ma X., Hopkins P. F., Faucher-Giguère C.-A., Zolman N., Muratov A. L., Kereš D., Quataert E., 2016a, MNRAS, 456, 2140
- Ma X., Hopkins P. F., Kasen D., Quataert E., Faucher-Giguère C.-A., Kereš D., Murray N., Strom A., 2016b, MNRAS, 459, 3614
- Ma X. et al., 2019, MNRAS, 487, 1844
- Pallottini A. et al., 2019, MNRAS, 487, 1689
- Pellegrini E. W., Reissl S., Rahner D., Klessen R. S., Glover S. C. O., Pakmor R., Herrera-Camus R., Grand R. J. J., 2019, preprint (arXiv:1905.04158)
- Pillepich A. et al., 2018, MNRAS, 475, 648
- Rosdahl J. et al., 2018, MNRAS, 479, 994
- Steidel C. C., Strom A. L., Pettini M., Rudie G. C., Reddy N. A., Trainor R. F., 2016, ApJ, 826, 159
- Torrey P. et al., 2018, MNRAS, 477, L16
- Torrey P. et al., 2019, MNRAS, 484, 5587
- Tremonti C. A. et al., 2004, ApJ, 613, 898
- Woosley S. E., Weaver T. A., 1995, ApJS, 101, 181
- Wu X., Davé R., Tacchella S., Lotz J., 2019, preprint (arXiv:1911.06330)
- Xiao L., Stanway E. R., Eldridge J. J., 2018, MNRAS, 477, 904

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The subcategories in boldface containing the word 'individual' are intended for use with specific astronomical objects; these should never be used alone, but always in combination with the most common names for the astronomical objects in question. Note that each object counts as one subcategory within the allowed limit of six.

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General

editorials, notices errata, addenda extraterrestrial intelligence history and philosophy of astronomy miscellaneous obituaries, biographies publications, bibliography sociology of astronomy standards

Physical data and processes

acceleration of particles accretion, accretion discs asteroseismology astrobiology astrochemistry astroparticle physics atomic data atomic processes black hole physics chaos conduction convection dense matter diffusion dynamo elementary particles equation of state gravitation gravitational lensing: micro gravitational lensing: strong gravitational lensing: weak gravitational waves hydrodynamics instabilities line: formation line: identification line: profiles magnetic fields magnetic reconnection (magnetohydrodynamics) MHD masers molecular data molecular processes neutrinos nuclear reactions, nucleosynthesis, abundances opacity plasmas polarization

radiation: dynamics radiation mechanisms:general radiation mechanisms: non-thermal radiation mechanisms: thermal radiative transfer relativistic processes scattering shock waves solid state: refractory solid state: volatile turbulence waves

Astronomical instrumentation, methods and techniques atmospheric effects balloons instrumentation: adaptive optics instrumentation: detectors instrumentation: high angular resolution instrumentation: interferometers instrumentation: miscellaneous instrumentation: photometers instrumentation: polarimeters instrumentation: spectrographs light pollution methods: analytical methods: data analysis methods: laboratory: atomic methods: laboratory: molecular methods: laboratory: solid state methods: miscellaneous methods: numerical methods: observational methods: statistical site testing space vehicles space vehicles: instruments techniques: high angular resolution techniques: image processing techniques: imaging spectroscopy techniques: interferometric techniques: miscellaneous techniques: photometric techniques: polarimetric techniques: radar astronomy techniques: radial velocities techniques: spectroscopic telescopes

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Software

software: data analysis software: development software: documentation software: public release software: simulations

Astrometry and celestial mechanics

astrometry celestial mechanics eclipses ephemerides occultations parallaxes proper motions reference systems time

The Sun

Sun: abundances Sun: activity Sun: atmosphere Sun: chromosphere Sun: corona Sun: coronal mass ejections (CMEs) Sun: evolution Sun: faculae, plages Sun: filaments, prominences Sun: flares Sun: fundamental parameters Sun: general Sun: granulation Sun: helioseismology Sun: heliosphere Sun: infrared Sun: interior Sun: magnetic fields Sun: oscillations Sun: particle emission Sun: photosphere Sun: radio radiation Sun: rotation (Sun:) solar-terrestrial relations (Sun:) solar wind (Sun:) sunspots Sun: transition region Sun: UV radiation Sun: X-rays, gamma-rays

Planetary systems comets: general

comets: individual: ... Earth interplanetary medium Kuiper belt: general

Kuiper belt objects: individual: ... meteorites, meteors, meteoroids minor planets, asteroids: general

minor planets, asteroids: individual: . . .

Moon Oort Cloud planets and satellites: atmospheres planets and satellites: aurorae planets and satellites: composition planets and satellites: detection planets and satellites: dynamical evolution and stability planets and satellites: formation planets and satellites: fundamental parameters planets and satellites: gaseous planets planets and satellites: general

planets and satellites: individual: ...

planets and satellites: interiors planets and satellites: magnetic fields planets and satellites: oceans planets and satellites: physical evolution planets and satellites: rings planets and satellites: surfaces planets and satellites: tectonics planets and satellites: terrestrial planets planet-disc interactions planet-star interactions protoplanetary discs zodiacal dust

Stars

stars: abundances stars: activity stars: AGB and post-AGB stars: atmospheres (stars:) binaries (including multiple): close (stars:) binaries: eclipsing (stars:) binaries: general (stars:) binaries: spectroscopic (stars:) binaries: symbiotic (stars:) binaries: visual stars: black holes (stars:) blue stragglers (stars:) brown dwarfs stars: carbon stars: chemically peculiar stars: chromospheres (stars:) circumstellar matter stars: coronae stars: distances stars: dwarf novae stars: early-type stars: emission-line, Be stars: evolution stars: flare stars: formation stars: fundamental parameters (stars:) gamma-ray burst: general (stars:) gamma-ray burst: individual: ... stars: general (stars:) Hertzsprung-Russell and colour-magnitude diagrams stars: horizontal branch stars: imaging stars: individual: ... stars: interiors

stars: jets stars: kinematics and dynamics stars: late-type stars: low-mass stars: luminosity function, mass function stars: magnetars stars: magnetic field stars: massive stars: mass-loss stars: neutron (stars:) novae, cataclysmic variables stars: oscillations (including pulsations) stars: peculiar (except chemically peculiar) (stars:) planetary systems stars: Population II stars: Population III stars: pre-main-sequence stars: protostars (stars:) pulsars: general (stars:) pulsars: individual: ... stars: rotation stars: solar-type (stars:) starspots stars: statistics (stars:) subdwarfs (stars:) supergiants (stars:) supernovae: general (stars:) supernovae: individual: ... stars: variables: Cepheids stars: variables: Scuti stars: variables: general stars: variables: RR Lyrae stars: variables: S Doradus stars: variables: T Tauri, Herbig Ae/Be (stars:) white dwarfs stars: winds, outflows stars: Wolf-Rayet

Interstellar medium (ISM), nebulae

ISM: abundances ISM: atoms ISM: bubbles ISM: clouds (*ISM:*) cosmic rays (*ISM:*) dust, extinction ISM: evolution ISM: general (*ISM:*) HII regions (*ISM:*) Herbig–Haro objects

ISM: individual objects: ...

(except planetary nebulae) ISM: jets and outflows ISM: kinematics and dynamics ISM: lines and bands ISM: magnetic fields ISM: molecules (ISM:) photodissociation region (PDR) (ISM:) planetary nebulae: general (ISM:) planetary nebulae: individual: ... ISM: structure ISM: supernova remnants

The Galaxy

Galaxy: abundances Galaxy: bulge Galaxy: centre Galaxy: disc Galaxy: evolution Galaxy: formation Galaxy: fundamental parameters Galaxy: general (Galaxy:) globular clusters: general (Galaxy:) globular clusters: individual: ... Galaxy: halo Galaxy: kinematics and dynamics (Galaxy:) local interstellar matter Galaxy: nucleus (Galaxy:) open clusters and associations: general (Galaxy:) open clusters and associations: individual: ... (Galaxy:) solar neighbourhood Galaxy: stellar content Galaxy: structure

Galaxies

galaxies: abundances galaxies: active galaxies: bar (galaxies:) BL Lacertae objects: general (galaxies:) **BL Lacertae objects: individual:...** galaxies: bulges galaxies: clusters: general

galaxies: clusters: individual: ...

galaxies: clusters: intracluster medium galaxies: disc galaxies: distances and redshifts galaxies: dwarf galaxies: elliptical and lenticular, cD galaxies: evolution galaxies: formation galaxies: fundamental parameters galaxies: general galaxies: groups: general

galaxies: groups: individual: ...

galaxies: haloes galaxies: high-redshift

galaxies: individual: . . .

galaxies: interactions (galaxies:) intergalactic medium galaxies: irregular galaxies: ISM galaxies: jets galaxies: kinematics and dynamics (galaxies:) Local Group galaxies: luminosity function, mass function (galaxies:) Magellanic Clouds galaxies: magnetic fields galaxies: nuclei galaxies: peculiar galaxies: photometry (galaxies:) quasars: absorption lines (galaxies:) quasars: emission lines (galaxies:) quasars: general

(galaxies:) quasars: individual: ...

(galaxies:) quasars: supermassive black holes galaxies: Seyfert galaxies: spiral galaxies: starburst galaxies: star clusters: general

galaxies: star clusters: individual: ...

galaxies: star formation galaxies: statistics galaxies: stellar content galaxies: structure

Cosmology

(cosmology:) cosmic background radiation (cosmology:) cosmological parameters (cosmology:) dark ages, reionization, first stars (cosmology:) dark energy (cosmology:) dark matter (cosmology:) diffuse radiation (cosmology:) distance scale (cosmology:) early Universe (cosmology:) inflation (cosmology:) large-scale structure of Universe cosmology: miscellaneous cosmology: observations (cosmology:) primordial nucleosynthesis cosmology: theory

Resolved and unresolved sources as a function of wavelength

gamma-rays: diffuse background gamma-rays: galaxies gamma-rays: galaxies: clusters gamma-rays: general gamma-rays: ISM gamma-rays: stars infrared: diffuse background infrared: galaxies infrared: general infrared: ISM infrared: planetary systems infrared: stars radio continuum: galaxies radio continuum: general radio continuum: ISM radio continuum: planetary systems radio continuum: stars radio continuum: transients radio lines: galaxies radio lines: general radio lines: ISM radio lines: planetary systems radio lines: stars submillimetre: diffuse background submillimetre: galaxies submillimetre: general submillimetre: ISM submillimetre: planetary systems submillimetre: stars ultraviolet: galaxies

ultraviolet: general ultraviolet: ISM ultraviolet: planetary systems ultraviolet: stars X-rays: binaries X-rays: bursts X-rays: diffuse background X-rays: galaxies X-rays: galaxies: clusters X-rays: general X-rays: individual: ... X-rays: ISM X-rays: stars

Transients

(transients:) black hole mergers (transients:) black hole - neutron star mergers (transients:) fast radio bursts (transients:) gamma-ray bursts (transients:) neutron star mergers transients: novae transients: supernovae transients: tidal disruption events